

## Structure, Dielectric Properties and Impedance Spectroscopy of BaTi<sub>0.97</sub>(Zn,V,Ta)<sub>0.03</sub>O<sub>3</sub> Ceramics

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**Summary:** BaTi<sub>0.97</sub>Zn<sub>0.01</sub>V<sub>0.01</sub>Ta<sub>0.01</sub>O<sub>3</sub> sample was prepared using the conventional mixed oxide sintering route. The samples were characterized using X-ray diffraction, Raman spectroscopy, Scanning Electron Microscopy and Impedance spectroscopy. Structure of the sample was analyzed using Rietveld analysis which showed the formation of tetragonal symmetry. Raman spectroscopy was also in agreement with the XRD results, further indicating the ferroelectric signature in the sample. Electrical microstructure analyzed by impedance spectroscopy analysis revealed two electroactive regions in the sample, belonging to the grain boundary and bulk. An extrinsic conduction mechanism was observed across the sample. Dielectric properties as a function of temperature demonstrated phase transition ( $T_C$ ) at  $\sim 94$  °C, with relative permittivity  $\sim 3280$  at  $T_C$ . Antiferroelectric-like double hysteresis loops were observed in P-E analysis. The energy storage density was calculated to be 0.26 J/cm<sup>3</sup> at an electric field of 50 kV/cm, indicating that BaTi<sub>0.97</sub>Zn<sub>0.01</sub>V<sub>0.01</sub>Ta<sub>0.01</sub>O<sub>3</sub> ceramics can be potentially useful base materials for high density energy storage capacitors.

**Keywords:** BaTiO<sub>3</sub>; Ferroelectrics; Dielectric properties; Impedance spectroscopy, Ceramics.

### Introduction

The importance of a dielectric material can be determined primarily by its relative permittivity, temperature stability, and breakdown strength for use in capacitors and high energy density storage applications. BaTiO<sub>3</sub> (BT) is an important perovskite structured material. BT-based ceramics are mostly used in MLCC (multi-layer ceramic capacitors), PTCR (positive temperature coefficient resistors) and thermistors [1, 2]. At Curie temperature ( $T_C$ ), relative permittivity ( $\epsilon_r$ ) of BT is  $\sim 10000$  [3] that varies with grain size [4]. The optimum sintering temperature (ST) of BT ranges from 1350-1400°C and different methods are used to lower ST of BT based ceramics [5-8]. It is technologically as well as commercially important to lower ST to cut down the processing costs and to lower it than the melting point of electrode such as Ni and Cu in MLCCs. Ceramic capacitor can store electric energy with much higher power output than the conventional batteries. However, their poor storage capacity compared to batteries is a challenge for researchers. Anti-ferroelectric materials have far higher power density than ferroelectrics. So far, lead based materials have shown excellent properties because of their anti-ferroelectric characters. Lead free ceramics are preferred because of the toxic nature of lead (Pb) and BT-based materials are one of the choices because of its high  $\epsilon_r$ . The co-doping of Ta<sub>2</sub>O<sub>5</sub> and ZnO in BT has been reported to increase  $T_c$  and  $\epsilon_r$  along with a

decrease in sintering temperature [9]. The melting point of Zn and V is low which may be useful to decrease the sintering temperature in comparison to pure BT. In this study, Zn, Ta and V were substituted for Ti in BT according to the formula BaTi<sub>0.97</sub>Zn<sub>0.01</sub>V<sub>0.01</sub>Ta<sub>0.01</sub>O<sub>3</sub> and the resulting ceramics and their structure and electrical properties were studied.

### Experimental

BaTi<sub>0.97</sub>Zn<sub>0.01</sub>V<sub>0.01</sub>Ta<sub>0.01</sub>O<sub>3</sub> samples were prepared using reagent grade BaCO<sub>3</sub>, TiO<sub>2</sub>, ZnO, V<sub>2</sub>O<sub>5</sub> and Ta<sub>2</sub>O<sub>5</sub> (Sigma Aldrich) through solid state sintering route. All reagents were dried at appropriate temperatures to get rid of the adsorbed moisture. The dried reagents were batched according to the molar ratios of the composition and then ball milled for 12 h. After ball milling, the powders were calcined at 950 °C for 6 h at a ramp rate of 5 °C/min. The calcined powders were re-milled again to reduce particle size and get rid of any agglomerates. The resulting free flowing powders were pressed into green pellets using a uniaxial pellet presser at  $\sim 50$  MPa. The green pellets were densified into compact ceramics by sintering at 1100-1250 °C for 4 h at a heating/cooling rate of 5 °C/min.

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The density of sintered samples was measured using a high precision densitometer (MDs-300). Phase of the sintered samples was identified using Siemens X-ray diffractometer with Cu-K $\alpha$  radiation source. Rietveld analysis of the sample was carried out using GSAS+EXPGUI computer program [10]. Raman spectra of the samples were obtained using a Renishaw Raman spectrometer (New Mills, Wotton-Under-Eagle, UK). Polarization vs electric field (P-E) loops were obtained using a ferroelectric test system. A JSM-5910 (JEOL) SEM (scanning electron microscope) was used for microstructural analysis. Temperature dependent  $\epsilon_r$  and dielectric loss (tan $\delta$ ) were studied in the temperature range ~ 20 to 250 °C using an Agilent LCR meter (HP 4284A). Complex Impedance spectroscopy (CIS) of the samples was performed using an Agilent (E4980A) Impedance Analyzer in the temperature range 300-500 °C.

## Results and Discussion

X-rays diffraction trace recorded for sintered BaTi<sub>0.97</sub>V<sub>0.01</sub>Ta<sub>0.01</sub>Zn<sub>0.01</sub>O<sub>3</sub> sample is shown in Fig. 1. The XRD peaks matched PDF # 74-1959 for tetragonal BT with space group *P4mm*. A low intensity additional peak (labeled as '\*') was observed at ~ 27.6° which could not be identified. The intensity ratio of the observed (002) and (200) XRD peaks was 1:2 which further confirmed tetragonal symmetry for the sintered phase [11]. The strongest peak at 45.3° was selected and yielded the average crystallite size to be 35 nm. The structure was further confirmed using Rietveld analysis which converged for the tetragonal (*P4mm*) structure. The refined lattice parameters were  $a = b = 4.0003(2)$  Å and  $c = 4.0226(3)$  Å. The extracted refined parameters are listed in Table-1. Here, Ba<sup>2+</sup> with larger ionic radius occupy A-site and intermediate size cations (Ti<sup>4+</sup>, Ta<sup>5+</sup>, Zn<sup>2+</sup> and V<sup>5+</sup>) occupy B-site of the host lattice. Oxygen (O1 and O3) ions are at the anionic site

occupying 1*b* and 2*c* positions, respectively as shown in Fig. 2.

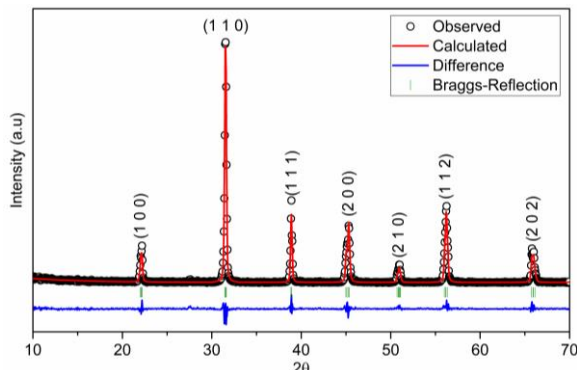


Fig. 1: Observed, calculated and difference in the XRD patterns of BaTi<sub>0.97</sub>(Zn,V,Ta)<sub>0.03</sub>O<sub>3</sub> sample sintered at 1200°C for 4h.

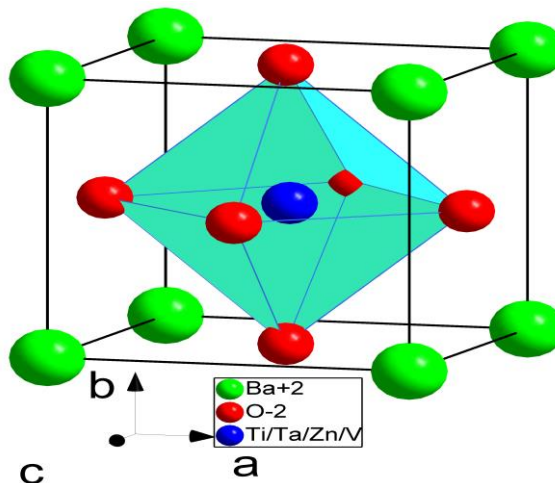


Fig. 2: Unit cell of BaTi<sub>0.97</sub>(Zn,V,Ta)<sub>0.03</sub>O<sub>3</sub>.

Table-1: Structural data of BaTi<sub>0.97</sub>V<sub>0.01</sub>Ta<sub>0.01</sub>Zn<sub>0.01</sub>O<sub>3</sub> sample.

Crystal system	Tetragonal					
Space group	<i>P4mm</i> (no. 99)					
Unit cell dimensions	$a = 4.0003(2)$ Å $c = 4.0226(3)$ Å					
Cell volume	64.37(0) Å <sup>3</sup>					
Z	1					
Density, calculated	6.064 g/cm <sup>3</sup>					
$R_p$	0.1454					
$R_{wp}$	0.2092					
Goodness of fit ( $\chi^2$ )	1.45					
Atomic coordinates and isotropic displacement parameters						
Atom	Wyckoff Position	Occupancy	x	y	z	U
Ba <sup>2+</sup>	1 <i>a</i>	1	0	0	0	0.0071
Ti <sup>4+</sup>	1 <i>b</i>	0.97	1/2	1/2	0.48200	0.0112
Ta <sup>5+</sup> /V <sup>5+</sup> /Zn <sup>2+</sup>	1 <i>b</i>	0.01	1/2	1/2	0.48200	0.0112
O1 (O <sup>2-</sup> )	1 <i>b</i>	1	1/2	1/2	-0.04582	0.0000
O2 (O <sup>2-</sup> )	2 <i>c</i>	1	1/2	0	0.51451	0.0000

The crystallite size was determined using the full widths at half maximum (FWHM) employing the Debye–Scherrer equation (1)

$$d = \frac{k\lambda}{\beta\cos\theta} \quad (1)$$

where  $d$  is the crystallite size, and  $\lambda$  is the wavelength of Cu  $K\alpha$ -X-ray radiation (0.15405 nm),  $\beta$  is the full-width half maxima (FWHM) in radians and  $\theta$  is the diffraction angle. The secondary electron SEM image of  $\text{BaTi}_{0.97}\text{V}_{0.01}\text{Ta}_{0.01}\text{Zn}_{0.01}\text{O}_3$  sample sintered at 1200°C for 4h is shown in Fig. 3. The dimensions of the observed grains of almost uniform contrast varied from ~1 to 4  $\mu\text{m}$ .

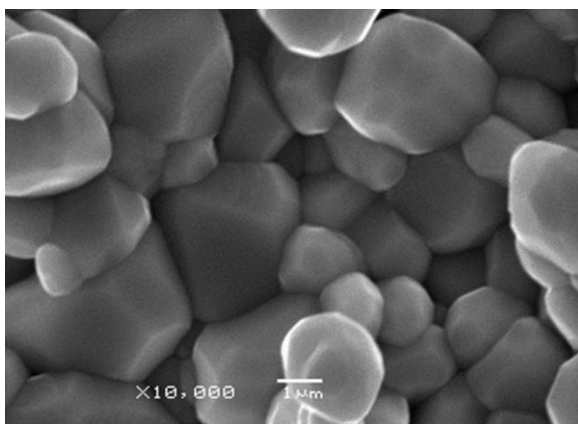


Fig. 3: SEM micrograph of  $\text{BaTi}_{0.97}(\text{Zn},\text{V},\text{Ta})_{0.03}\text{O}_3$  sample sintered at 1200°C for 4h.

The Raman spectrum of  $\text{BaTi}_{0.97}\text{V}_{0.01}\text{Ta}_{0.01}\text{Zn}_{0.01}\text{O}_3$  sample taken at room temperature (RT) is shown in Fig. 4.  $\text{BaTiO}_3$  has five atoms and fifteen degrees of freedom per unit cell [12]. The  $180\text{ cm}^{-1}$  mode is ascribed to interference of Raman scattering between two vibrational modes with overlapping frequency. The interference arises due to an an-harmonic band between the two layer vibrational modes [13]. Tetragonal BT phase is known to exhibit Raman scattering bands at ~250, 515 and  $717\text{ cm}^{-1}$  and a sharp peak at ~306  $\text{cm}^{-1}$ . Theoretically, the phase with cubic ( $Pm\bar{3}m$ ) symmetry exhibits no Raman active common modes but commonly exhibits broad bands at ~250 and  $515\text{ cm}^{-1}$ , originating from the local-disorder coupled with the position of  $\text{Ti}^{4+}$  ions. The bands near  $717\text{ cm}^{-1}$  in the tetragonal structure is due to  $\text{Ba}^{2+}$  defects in the BT lattice. A small mode appearing at ~838  $\text{cm}^{-1}$  is due to the aliovalent substituents on the B-site in BT while such a mode does not appear due to substitutions on the A-site in BT [14]. This mode arises due to the charge difference at equivalent sites in BT. In conclusion, the Raman spectrum confirms that the V-, Ta- and Zn-doped composition

crystallized into a tetragonal symmetry, consistent with the previous studies [15-17].

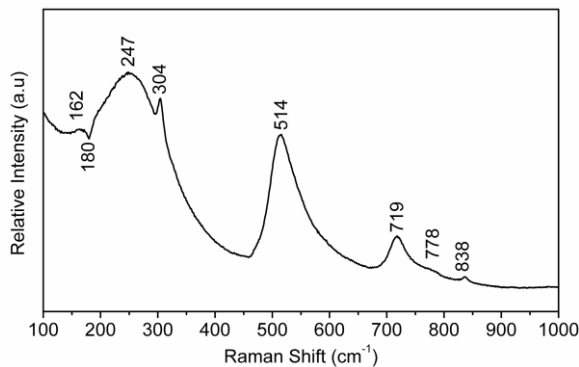


Fig. 4: Raman spectrum of  $\text{BaTi}_{0.97}(\text{Zn},\text{V},\text{Ta})_{0.03}\text{O}_3$  ceramics.

To analyze the electrical microstructure of the sample, impedance spectroscopy which is considered as a powerful tool to differentiate between different electroactive regions? The complex plane plots ( $Z''$  versus  $Z'$ ) at 300 and 450°C showed a large semicircle at low frequency and smaller semicircle at high frequency (Fig. 5) which indicated the presence of two different conduction mechanisms. The large semicircle observed at low frequency was indicative of a relatively more resistive region while the small semicircle observed at high frequency was indicative of a relatively more conductive region. The arc at high frequency is related to the bulk while the arc at low frequency originates from contributions due to the grain boundaries. The modulus plot also showed the same behavior, Fig. 6.

To further confirm, the combined spectroscopic plot ( $Z''$  and  $M''$  versus  $\log f$ ) was also analyzed, Fig. 7. Consistent with the complex plane plot (Fig. 5), two peaks were observed on the spectroscopic plot ( $Z''$  versus  $\log f$ ) as well. The value of resistance (R) and capacitance (C) were obtained from the peak maxima in the spectroscopic plots using formula  $R = 2Z''_{\text{max}}$  and  $C = 1/(2\pi f_{\text{max}}R)$ , respectively. At 450 °C, the capacitance and resistance values calculated from  $Z''$  versus  $\log f$  plot were  $1.3 \times 10^{-9}\text{ F}$  and  $1.5\text{ M}\Omega$ , respectively. Similarly, the capacitance and resistance values obtained from the  $M''$  versus  $\log f$  plot were  $27.2 \times 10^{-12}\text{ F}$  and  $5841\ \Omega$ , respectively. The capacitance versus  $\log f$  plot shown in Fig. 8 was found in good agreement with the capacitance values associated with the bulk ( $10^{-12}\text{ F}$ ) and grain boundary ( $10^{-9}\text{ F}$ ) [18, 19]. To study the conduction mechanism in the sample, activation energies were calculated from the Arrhenius plot of conductivity (Fig. 9), using equation (2)

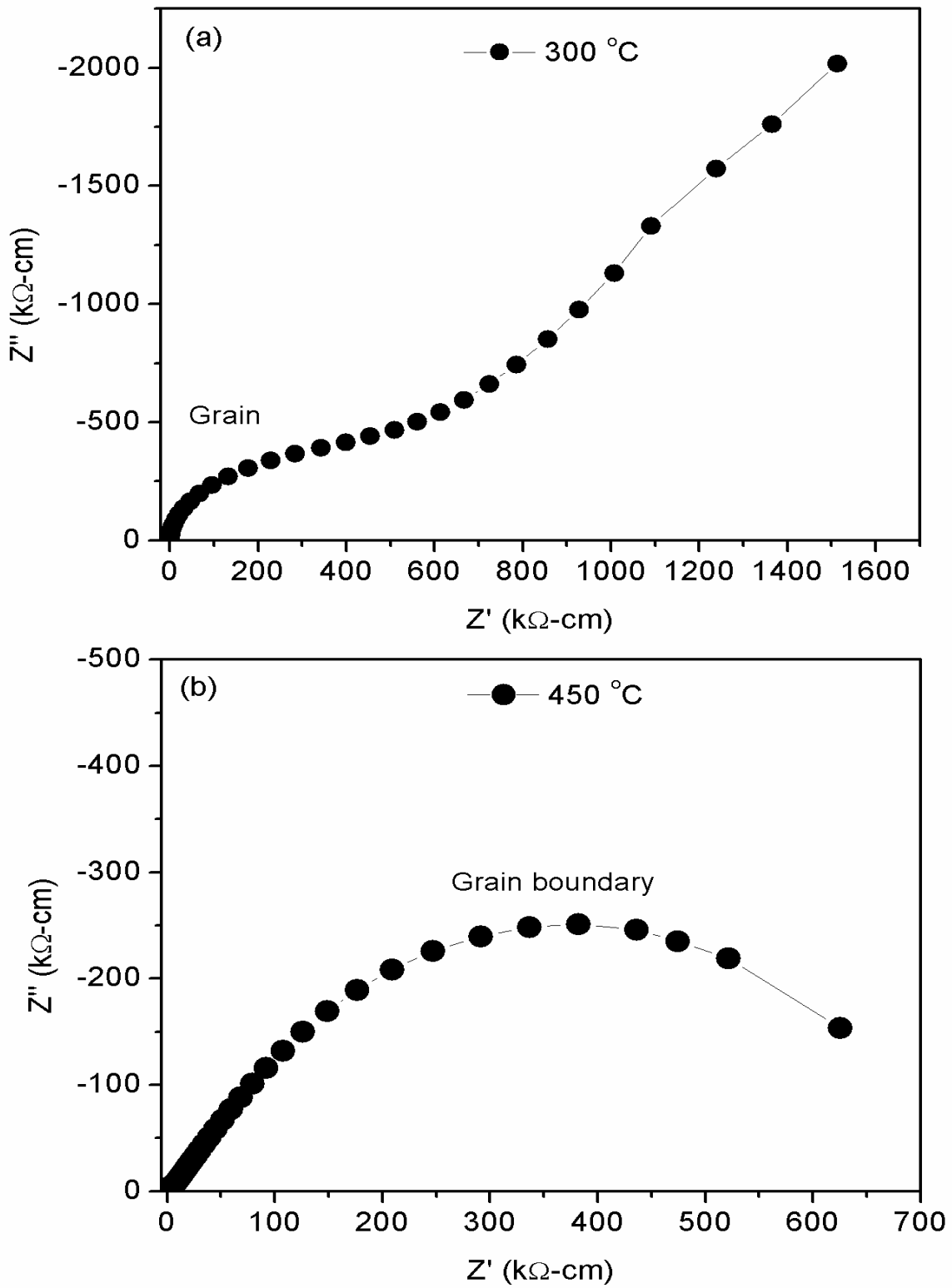


Fig. 5: Complex plane plots ( $Z'$  vs  $Z''$ ) at a) 300°C and b) 450°C.

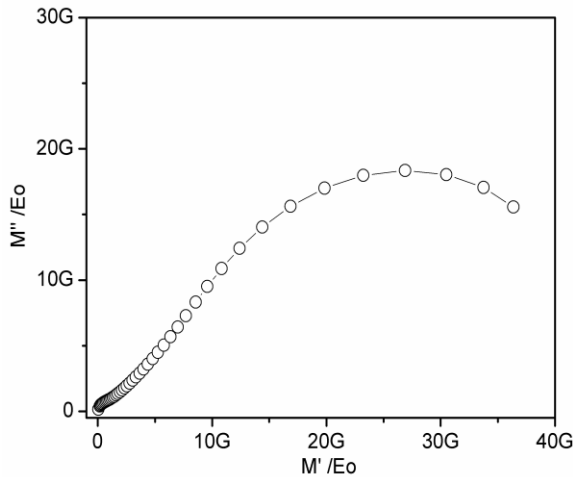


Fig. 6: Modulus plot ( $M'$  vs  $M''$ ) at 450°C.

$$\square = \square_0 \exp\left(\frac{-E_a}{k_B T}\right) \quad (2)$$

where  $\square_0$  is the pre-exponential factor,  $T$  is the temperature,  $E_a$  is activation energy and  $k_B$  is Boltzmann constant.  $E_a$  of the bulk and grain boundary were calculated to be 1.04 eV and 1.22 eV, respectively, which suggest that the extrinsic conduction (oxygen ion conduction) takes place in the samples which is commonly reported for pure BT.

$\epsilon_r$  and  $\tan\delta$  plots of  $\text{BaTi}_{0.97}\text{Zn}_{0.01}\text{V}_{0.01}\text{Ta}_{0.01}\text{O}_3$  at 1 kHz to 1 MHz in the temperature range ~25-250 °C are shown in Fig. 10.  $\epsilon_r$  at the peak temperature was observed to be ~ 3280. The phase transition temperature (Curie temperature  $T_C$ ) of undoped BT is commonly reported to be ~ 125 °C; however, doping with V, Ta and Zn shifted the Curie temperature to a relatively lower temperature (~ 94 °C). The Curie temperature of BT basically depends on the domain structure, particle size and electrical behavior of the grain boundary [20-22] which makes the explanation of the phase transition behavior of BT difficult. Similarly, covalency also affects  $\epsilon_r$  of BT which may be a possible reason for the observed decrease in  $T_C$  [20, 23].  $\tan\delta$  was ~ 0.008 at room temperature.

The P-E loops of  $\text{BaTi}_{0.97}\text{Zn}_{0.01}\text{V}_{0.01}\text{Ta}_{0.01}\text{O}_3$  sample were measured by exposing the samples to electric fields between 10 - 50 kV/cm as shown in Fig. 11. The sample exhibited antiferroelectric like behavior at all E values showing nearly zero

polarization at room temperature, suggesting that the total energy can be released by removing electric field. The observed double hysteresis loop is typical of antiferroelectric components prototyped by  $\text{PbZrO}_3$  [24]. It should be noted that double hysteresis can have other origins such as Paraelectric-Ferroelectric (PE-FE) phase transition near  $T_C$  and / or an aging effect well below  $T_C$  [25]. Since, the observation of double hysteresis loops in the present study were at room temperature (well below  $T_C$ ); therefore, these could not be attributed to the PE-FE phase transition but are more likely to be arising of an effect similar to aging, as observed by Ren et al for aged Mn doped BT [26]. Generally double loops due to aging are attributed to gradual stabilization of domain patterns by defects such as vacancies, dopant and or impurities. Though these samples were not aged but the role of defects in pinching of the P-E loops cannot be excluded. Defects in the studied sample could be associated to presence of impurities such as  $\text{Fe}^{3+}$  in reagents and/or volatilization of  $\text{V}^{5+}$  during high temperature synthesis. The exact role of defects is not clear at this stage and warrants a systematic investigation, however regardless of the origin, double loops are favored for energy storage applications due to their low remnant polarization. Generally, KBT based ceramics have been extensively investigated for high energy density applications. More recently, Yu et al. reported 1.15 J/cm<sup>3</sup> under 105 kV/cm for BNT-BKT-BA system [27]. At  $E = 50$  kV/cm, the energy storage density was 0.26 J/cm<sup>3</sup> for  $\text{BaTi}_{0.97}\text{Zn}_{0.01}\text{V}_{0.01}\text{Ta}_{0.01}\text{O}_3$  sample which may be increased at higher E values, below breakdown of the sample. Therefore, further work is required to study the electric breakdown and hence maximum energy storage ability of  $\text{BaTi}_{0.97}\text{Zn}_{0.01}\text{V}_{0.01}\text{Ta}_{0.01}\text{O}_3$  sample. The energy efficiency of the sample at 50 kV/cm was found to be 72 %. There are several lead-based antiferroelectric materials that are found to be excellent for high energy storage applications, but environmentally friendly materials are required for these applications. On the other hand, there are very few lead-free antiferroelectric materials therefore,  $\text{BaTi}_{0.97}\text{Zn}_{0.01}\text{V}_{0.01}\text{Ta}_{0.01}\text{O}_3$  ceramics could be an ideal lead-free material for high energy storage applications. Furthermore, the sintering temperature of  $\text{BaTi}_{0.97}\text{Zn}_{0.01}\text{V}_{0.01}\text{Ta}_{0.01}\text{O}_3$  ceramics is almost 150 °C lower than the BT prepared via conventional mixed oxide sintering route [3].

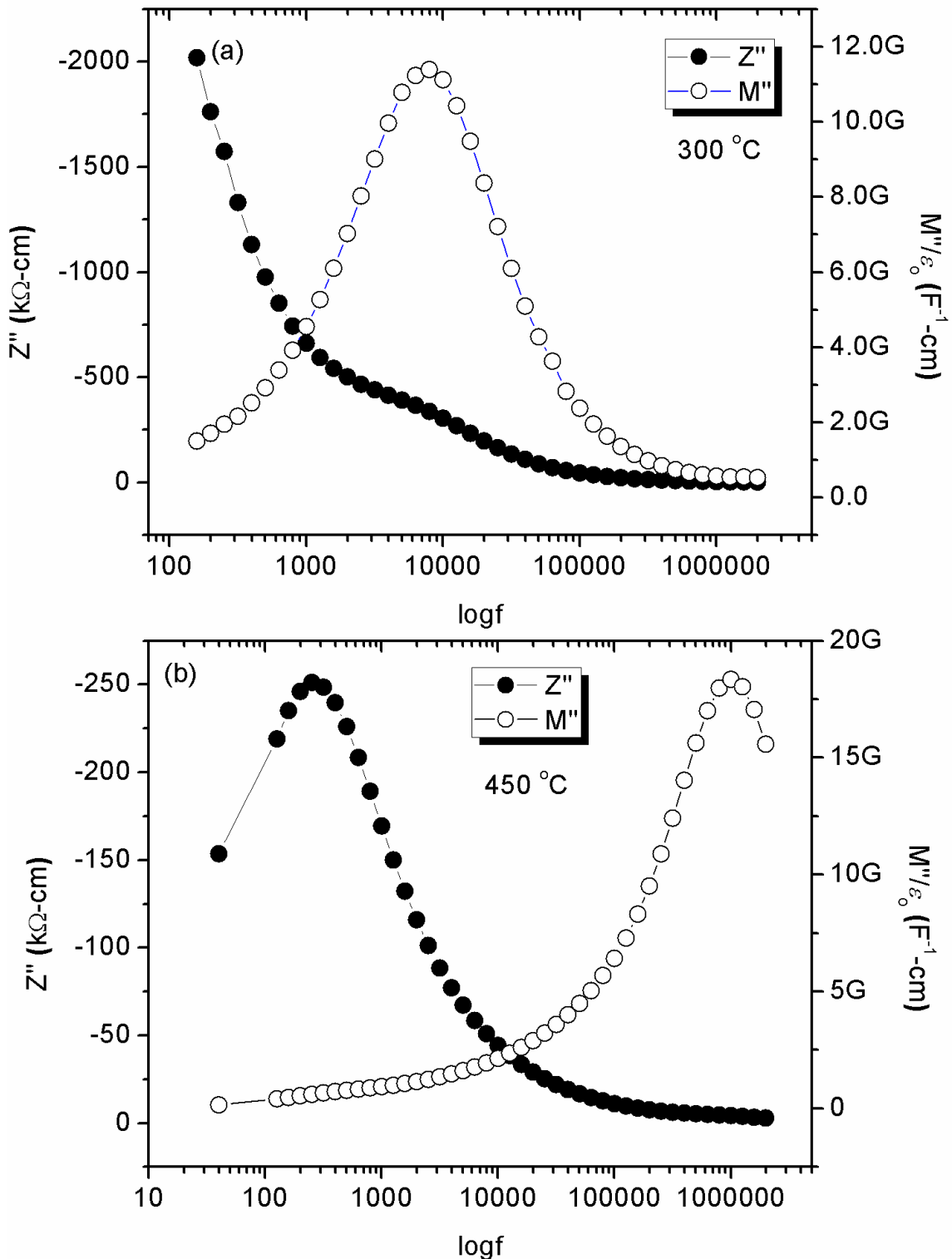


Fig. 7: Impedance spectroscopic plots of BaTi<sub>0.97</sub>(Zn,V,Ta)<sub>0.03</sub>O<sub>3</sub> ceramics at; a) 300°C and b) 450°C.

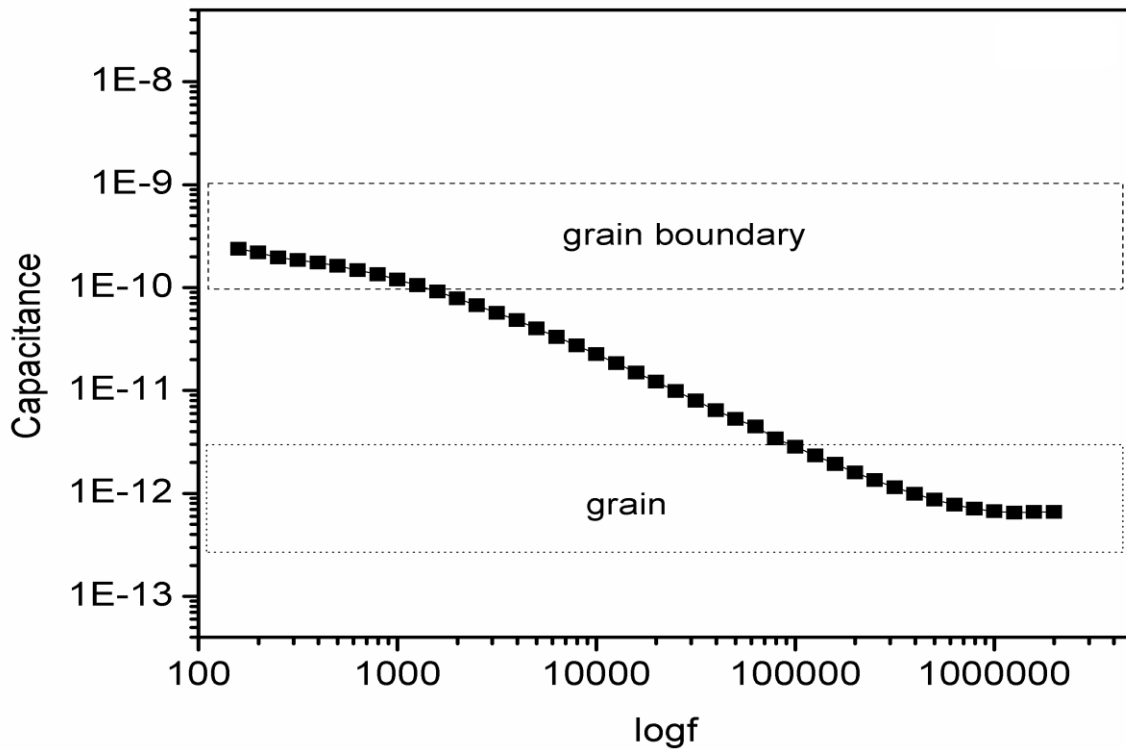


Fig. 8: Plot of logC vs. logf for BaTi<sub>0.97</sub>(Zn,V,Ta)<sub>0.03</sub>O<sub>3</sub> sample at 300°C.

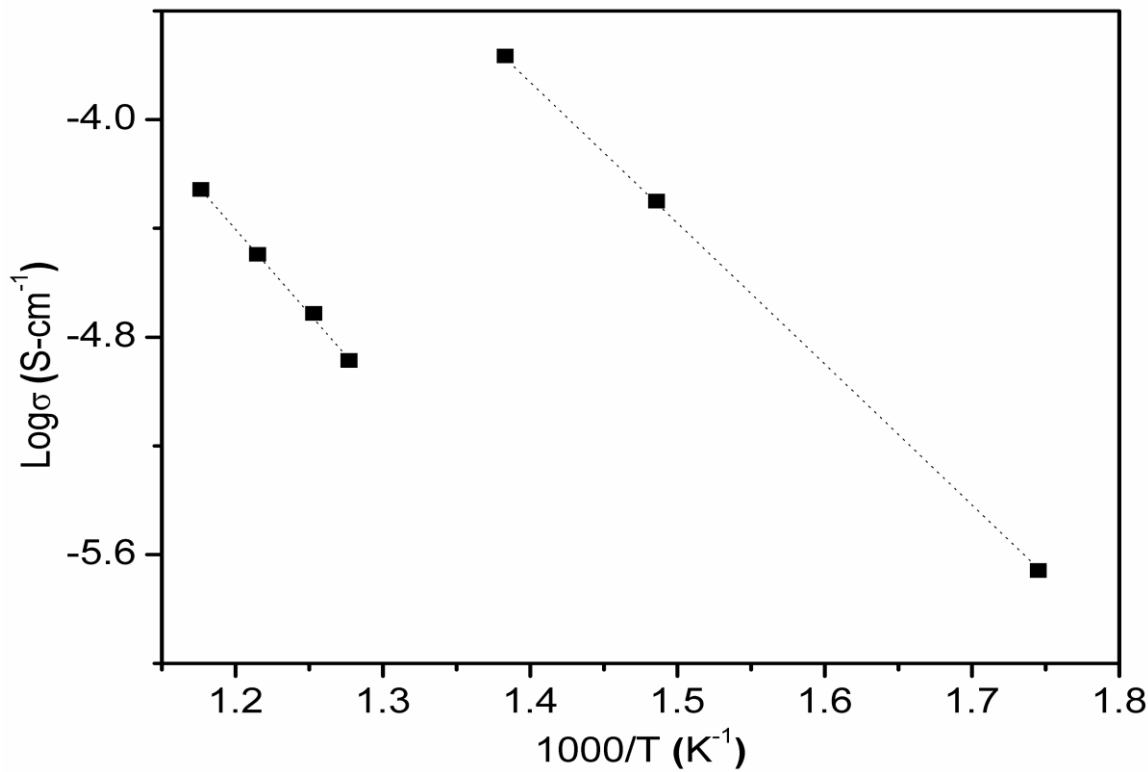


Fig. 9: Arrhenius plot of conductivity for BaTi<sub>0.97</sub>(Zn,V,Ta)<sub>0.03</sub>O<sub>3</sub> ceramics.

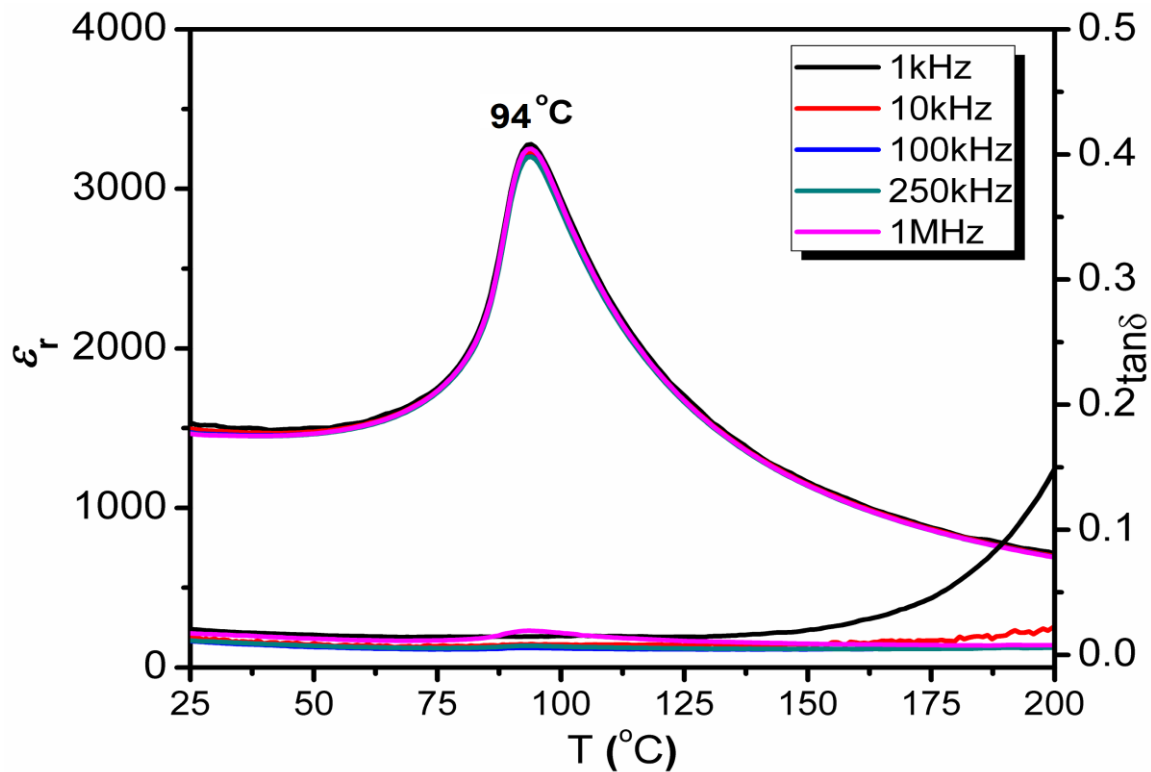


Fig. 10:  $\epsilon'$  and  $\tan\delta$  versus temperature for  $\text{BaTi}_{0.97}(\text{Zn,V,Ta})_{0.03}\text{O}_3$  at different frequencies.

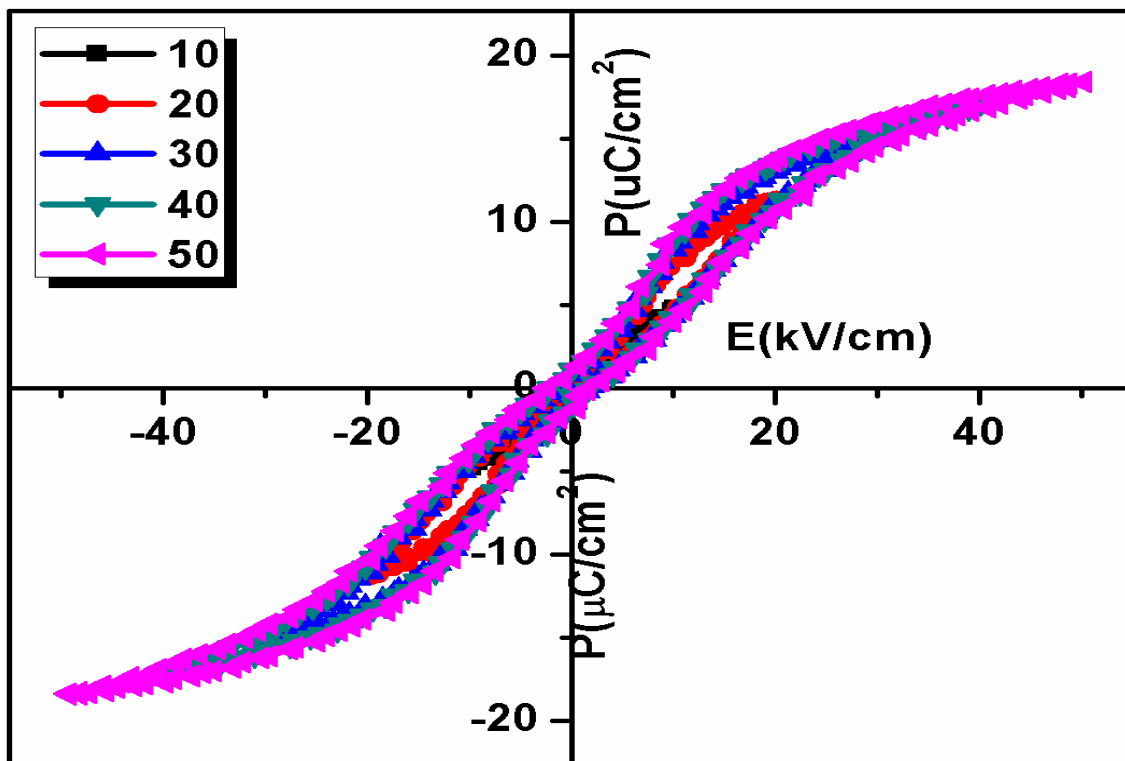


Fig. 11: P-E hysteresis loops of  $\text{BaTi}_{0.97}(\text{Zn,V,Ta})_{0.03}\text{O}_3$  ceramics.



## Conclusion

BaTi<sub>0.97</sub>Zn<sub>0.01</sub>V<sub>0.01</sub>Ta<sub>0.01</sub>O<sub>3</sub> samples were fabricated using conventional using mixed oxide solid sintering route. The dopants did not induce any changes in symmetry and the sample crystallized into a tetragonal symmetry, confirmed by XRD and Raman spectroscopy. Electrical microstructure of the sample was studied using impedance spectroscopy which revealed that extrinsic conduction take place in the sample. A double hysteresis loop was observed for the sample and the energy storage density was 0.26 J/cm<sup>3</sup> under 50 kV/cm showing potential of these ceramics to be used for high density energy storage capacitors.

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